# Particle classification 

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## Categories

Hadrons interact via strong force
Mesons spin 0,1 and $\mathrm{m}_{e}<\mathrm{m}<\mathrm{m}_{\mathrm{p}}$
Baryons spin $\frac{1}{2}, \frac{3}{2}, \frac{5}{2} m \geq m_{p}$. Including $p, n$. Decay products always include a $p$.

Leptons, spin $\frac{1}{2}$, appear to be structureless point particles.

- electron ( $e^{-}$), electron neutrino $v_{e}$
- tau $\tau^{-}$, tau neutrino $\gamma_{\tau}$
- mu, $\mu^{-}$, mu neutrino $\gamma_{\mu}$.

Plus their respective antiparticles.

## Neutrino physics

Neutrinos have a helicity: spin is aligned with $\overrightarrow{\mathbf{p}}$ for antineutrinos, opposite for neutrinos.

Neutrinos seem to have a small ( $\sim \mathrm{eV}$ ) mass. This is difficult to measure.

Neutrino oscillations have been recently confirmed. The different neutrion types can among themselves as they propagate from a source. This explains the solar neutrino problem and should permit determination of the neutrino masses.

## Conservation Laws

In nuclear reactions or decays, the following are conserved, adding up the numbers for the types of particles:

Baryon number $B= \pm 1$ for baryons/antibaryons, $B=0$ for others.

Lepton number $\mathrm{L}= \pm 1$ for electrons and their neutrinos, muon and tau families hav $\mathrm{L}=0$.

Strangeness Some heavy hadrons are produced in pairs in reactions, such as the $K, \Lambda, \Sigma$. They have strangeness numbers that are conserved in reactions.

Resonance particles

The very short lifetime $\Delta^{+}$, of mass $1231 \mathrm{MeV} / \mathrm{c}^{2}$. Consider

$$
e^{-}+p \rightarrow e^{-}+\Delta^{+}
$$

followed in $6 \times 10^{-24} \mathrm{~s}$ by

$$
\Delta^{+} \rightarrow \pi^{+}+\mathrm{n}
$$

and the direct reaction

$$
e^{-}+p \rightarrow e^{-}+\pi^{+}+n
$$

where no $\Delta^{+}$is produced. How to tell them apart? The $\Delta^{+}$will not last long enough to leave a track.

The decay of the $\Delta^{+}$must satisfy

$$
E_{\Delta}^{2}=\left(p_{\Delta} c\right)^{2}+\left(m_{\Delta} c^{2}\right)^{2}
$$

or

$$
m_{\Delta} c^{2}=\sqrt{E_{\Delta}^{2}-\left(p_{\Delta} c\right)^{2}}
$$

We can't measure $E_{\Delta}$ and $\overrightarrow{\mathbf{p}}_{\Delta}$, but after the decay we can measure the outgoing particle properties, so $\mathrm{E}_{\Delta}=\mathrm{E}_{\pi}+\mathrm{E}_{\mathrm{n}}$ and $\overrightarrow{\mathbf{p}}_{\Delta}=\overrightarrow{\mathbf{p}}_{\pi}+\overrightarrow{\mathbf{p}}_{\mathrm{n}}$, and

$$
m_{\Delta} c^{2}=\sqrt{\left(E_{\pi}+E_{n}\right)^{2}-\left(\overrightarrow{\mathbf{p}}_{\pi}+\overrightarrow{\mathbf{p}}_{n}\right)^{2} c^{2}}
$$

This will be 1231 MeV if a $\Delta^{+}$decay is involved. If not, a broad range of values is possible for the direct reaction.

A histogram of energy values for this type of reaction will show a peak at the particle energy for a rapidly decaying particle. See fig 15.8

