



## **PROBLEM:**

A plane-polarized electromagnetic wave of frequency  $\omega$  in free space is incident normally on the flat surface of a non-permeable medium of conductivity  $\sigma$  and dielectric constant  $\varepsilon$ .

(a) Calculate the amplitude and phase of the reflected wave relative to the incident wave for arbitrary  $\sigma$  and  $\varepsilon$ .

(b) Discuss the limiting case of a very poor and a very good conductor, and show that for a good conductor the reflection coefficient (ratio of reflected to incident intensity) is approximately

$$R \approx 1 - 2 \frac{\omega}{c} \delta$$

where  $\delta$  is the skin depth.

## **SOLUTION:**

(a) If we set up the problem in the usual way with the incident wave E traveling in the positive z direction, the transmitted wave E' traveling in the positive z direction, and the reflected wave E'' traveling in the negative z direction, we get from the boundary conditions:

$$E_0' = E_0 + E_0''$$
  
 $k' E_0' = k E_0 - k E_0''$ 

We can use both of these to solve for the reflected wave in terms of the incident wave:

$$E_{0}'' = \frac{k}{k'} E_{0} - \frac{k}{k'} E_{0}'' - E_{0}$$
$$E_{0}'' = \frac{k - k'}{k + k'} E_{0}$$

The wave number in free space is just  $k = \frac{\omega}{c}$ 

The wave number in the other material obeys:  $k = \sqrt{\epsilon(\omega)\mu_0}\omega$ 

For a conducting material we can use  $\epsilon(\omega) = \epsilon + i \frac{\sigma}{\omega}$  so that we have

$$k' = \sqrt{\epsilon \,\mu_0 + i \frac{\sigma}{\omega} \mu_0} \,\omega$$

Plugging these in above leads to:

$$\frac{E_0''}{E_0} = \frac{1 - \sqrt{\frac{\epsilon}{\epsilon_0} + i\frac{\sigma}{\epsilon_0\omega}}}{1 + \sqrt{\frac{\epsilon}{\epsilon_0} + i\frac{\sigma}{\epsilon_0\omega}}}$$

The problem asks us to explicitly expand this in terms of amplitude and phase. The amplitude of a complex number is defined as  $|z| = \sqrt{z^* z}$  so that we have

$$\left|\frac{E_{0}}{E_{0}}\right| = \sqrt{\left|\frac{1 - \sqrt{\frac{\epsilon}{\epsilon_{0}} - i\frac{\sigma}{\epsilon_{0}\omega}}}{1 + \sqrt{\frac{\epsilon}{\epsilon_{0}} - i\frac{\sigma}{\epsilon_{0}\omega}}}\right| \left|\frac{1 - \sqrt{\frac{\epsilon}{\epsilon_{0}} + i\frac{\sigma}{\epsilon_{0}\omega}}}{1 + \sqrt{\frac{\epsilon}{\epsilon_{0}} + i\frac{\sigma}{\epsilon_{0}\omega}}}\right|}$$
$$\left|\frac{E_{0}}{E_{0}}\right| = \sqrt{\frac{1 + \sqrt{\frac{\epsilon^{2}}{\epsilon_{0}^{2}} + \frac{\sigma^{2}}{\epsilon_{0}^{2}\omega^{2}}} - 2\Re\sqrt{\frac{\epsilon}{\epsilon_{0}} + i\frac{\sigma}{\epsilon_{0}\omega}}}{1 + \sqrt{\frac{\epsilon^{2}}{\epsilon_{0}^{2}} + \frac{\sigma^{2}}{\epsilon_{0}^{2}\omega^{2}}} + 2\Re\sqrt{\frac{\epsilon}{\epsilon_{0}} + i\frac{\sigma}{\epsilon_{0}\omega}}}\right|}$$

The phase of a complex number is  $Arg(z) = \tan^{-1}\left(-i\frac{z-z^*}{z+z^*}\right)$ 

$$Arg\left(\frac{E_0''}{E_0}\right) = \tan^{-1}\left(\frac{2\Im\sqrt{\frac{\epsilon}{\epsilon_0} - i\frac{\sigma}{\epsilon_0\omega}}}{1 - \sqrt{\frac{\epsilon^2}{\epsilon_0^2} + \frac{\sigma^2}{\epsilon_0^2\omega^2}}}\right)$$

(b) Discuss the limiting case of a very poor and a very good conductor, and show that for a good conductor the reflection coefficient (ratio of reflected to incident intensity) is approximately

$$R \approx 1 - 2 \frac{\omega}{c} \delta$$

where  $\delta$  is the skin depth.

For a very poor conductor  $\sigma \ll \varepsilon_0 \omega$  which means the same as  $\sigma/\varepsilon_0 \omega \ll 1$ . This leads to:

$$\left|\frac{E_0"}{E_0}\right| = \sqrt{\frac{1 + \frac{\epsilon}{\epsilon_0} - 2\sqrt{\frac{\epsilon}{\epsilon_0}}}{1 + \frac{\epsilon}{\epsilon_0} + 2\sqrt{\frac{\epsilon}{\epsilon_0}}}}$$
$$\left|\frac{E_0"}{E_0}\right| = \frac{1 - \sqrt{\frac{\epsilon}{\epsilon_0}}}{1 + \sqrt{\frac{\epsilon}{\epsilon_0}}}$$

Now recognize:

$$\Im \sqrt{x + iy} = \Im (x^2 + y^2)^{1/4} e^{\frac{i}{2} \tan^{-1}(y/x)}$$
$$\Im \sqrt{x + iy} = (x^2 + y^2)^{1/4} \sin (\frac{1}{2} \tan^{-1}(y/x))$$

If  $y \ll x$  then  $\tan^{-1}(y/x) = y/x$  and  $\sin(\frac{1}{2}y/x) = \frac{1}{2}y/x$  so that

$$\Im \sqrt{x+iy} = \frac{1}{2} \sqrt{x} y/x$$
 leading to:

$$Arg\left(\frac{E_0"}{E_0}\right) = \tan^{-1}\left(\frac{\sqrt{\frac{\epsilon}{\epsilon_0}}\frac{\sigma}{\epsilon\omega}}{\frac{\epsilon}{\epsilon_0}-1}\right)$$

We must be careful because the arctan leads to some ambiguity of what quadrant the phase should be in. We know from physical arguments that the phase difference should be  $\pi$  for a perfect dielectric, so that our answer should reduce to  $\pi$  for  $\sigma \rightarrow 0$ . The answer must be in the third quadrant:

$$Arg\left(\frac{E_0''}{E_0}\right) = \pi + \frac{\sqrt{\frac{\epsilon}{\epsilon_0}}}{\frac{\epsilon}{\epsilon_0} - 1} \frac{\sigma}{\epsilon \omega}$$

For a very good conductor  $\sigma \gg \varepsilon_0 \omega$ 

$$\left|\frac{E_0"}{E_0}\right| = \sqrt{\frac{1 + \frac{\sigma}{\epsilon_0 \omega} - \sqrt{2}\sqrt{\frac{\sigma}{\epsilon_0 \omega}}}{1 + \frac{\sigma}{\epsilon_0 \omega} + \sqrt{2}\sqrt{\frac{\sigma}{\epsilon_0 \omega}}}} \quad \text{which leads to} \quad \left|\frac{E_0"}{E_0}\right| = \sqrt{\frac{1 - \sqrt{2}\sqrt{\frac{\epsilon_0 \omega}{\sigma}}}{1 + \sqrt{2}\sqrt{\frac{\epsilon_0 \omega}{\sigma}}}} \quad \text{after dropping small terms}$$

Use the binomial expansion  $\sqrt{1+x} = 1 + (1/2)x + ...$  on top and bottom and drop all small terms

$$\left|\frac{E_0"}{E_0}\right| = \frac{1 - \frac{1}{2}\sqrt{2}\sqrt{\frac{\epsilon_0\omega}{\sigma}}}{1 + \frac{1}{2}\sqrt{2}\sqrt{\frac{\epsilon_0\omega}{\sigma}}} \quad \text{which leads to} \quad \left|\frac{E_0"}{E_0}\right| = \frac{1 + \frac{1}{2}\frac{\epsilon_0\omega}{\sigma} - \sqrt{2}\sqrt{\frac{\epsilon_0\omega}{\sigma}}}{1 - \frac{1}{2}\frac{\epsilon_0\omega}{\sigma}}$$
$$\left|\frac{E_0"}{E_0}\right| = 1 - \sqrt{2}\sqrt{\frac{\epsilon_0\omega}{\sigma}}$$

We can now calculate the reflection coefficient:

$$R = \left| \frac{E_0}{E_0} \right|^2 = \left( 1 - \sqrt{2} \sqrt{\frac{\epsilon_0 \omega}{\sigma}} \right)^2 = 1 + 2 \frac{\epsilon_0 \omega}{\sigma} - 2 \sqrt{2} \sqrt{\frac{\epsilon_0 \omega}{\sigma}}$$

Drop the smallest term

$$\boxed{R=1-2\sqrt{2}\sqrt{\frac{\epsilon_0\omega}{\sigma}}} \quad \text{or} \quad \boxed{R=1-2\frac{\omega}{c}\delta} \quad \text{where} \quad \delta=\sqrt{\frac{2}{\mu_0\sigma\omega}}$$

The phase for a good conductor can also be found:

$$Arg\left(\frac{E_0''}{E_0}\right) = \tan^{-1}\left(\frac{\sqrt{2}\sqrt{\frac{\sigma}{\epsilon_0\omega}}}{\frac{\sigma}{\epsilon_0\omega} - 1}\right)$$

Drop the 1 in the denominator because it is much smaller:

$$Arg\left(\frac{E_0''}{E_0}\right) = \tan^{-1}\left(\sqrt{2}\sqrt{\frac{\epsilon_0\omega}{\sigma}}\right)$$
$$Arg\left(\frac{E_0''}{E_0}\right) = \pi + \sqrt{\frac{2\epsilon_0\omega}{\sigma}}$$