



C. S. BAIRD

## Jackson 5.13 Homework Problem Solution

Dr. Christopher S. Baird  
University of Massachusetts Lowell



### **PROBLEM:**

A sphere of radius  $a$  carries a uniform surface-charge distribution  $\sigma$ . The sphere is rotated about a diameter with constant angular velocity  $\omega$ . Find the vector potential and magnetic-flux density both inside and outside the sphere.

### **SOLUTION:**

The current density in spherical coordinates is:

$$\mathbf{J} = \sigma \omega a \sin \theta \delta(r-a) \hat{\boldsymbol{\phi}}$$

$$\mathbf{A} = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} d\mathbf{x}'$$

$$\mathbf{A} = \frac{\mu_0}{4\pi} \int \frac{\sigma \omega a \sin \theta' \delta(r'-a) \hat{\boldsymbol{\phi}}'}{|\mathbf{x} - \mathbf{x}'|} d\mathbf{x}'$$

Expand the denominator in spherical harmonics.

$$\mathbf{A} = \frac{\mu_0}{4\pi} \int_0^{2\pi} \int_0^\pi \int_0^\infty \sigma \omega \delta(r'-a) \hat{\boldsymbol{\phi}}' 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r'^l}{r^{l+1}} Y_{lm}^*(\theta', \phi') Y_{lm}(\theta, \phi) r'^2 \sin^2 \theta' dr' d\theta' d\phi'$$

We must be careful and realize that  $\hat{\boldsymbol{\phi}}' \neq \hat{\boldsymbol{\phi}}$  because the primed variables are being integrated over. The best way to handle this is to use the expansion:  $\hat{\boldsymbol{\phi}}' = -\sin \phi' \hat{\mathbf{i}} + \cos \phi' \hat{\mathbf{j}}$

$$\mathbf{A} = \mu_0 \sigma \omega a \int_0^{2\pi} \int_0^\pi \int_0^\infty \delta(r'-a) (-\sin \phi' \hat{\mathbf{i}} + \cos \phi' \hat{\mathbf{j}}) \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r'^l}{r^{l+1}} Y_{lm}^*(\theta', \phi') Y_{lm}(\theta, \phi) \sin^2 \theta' dr' d\theta' d\phi'$$

for  $r < r'$  (inside the sphere).

$$\mathbf{A} = \mu_0 \sigma \omega a \int_0^{2\pi} \int_0^\pi \int_0^\infty \delta(r'-a) (-\sin \phi' \hat{\mathbf{i}} + \cos \phi' \hat{\mathbf{j}}) \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r'^{l+2}}{r^{l+1}} Y_{lm}^*(\theta', \phi') Y_{lm}(\theta, \phi) \sin^2 \theta' dr' d\theta' d\phi'$$

for  $r > r'$  (outside the sphere).

Now evaluate the delta functions and rearrange:

$$\mathbf{A} = \mu_0 \sigma \omega \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r^l}{a^{l-2}} Y_{lm}(\theta, \phi) I_{l,m} \quad (\text{inside})$$

$$\mathbf{A} = \mu_0 \sigma \omega \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{a^{l+3}}{r^{l+1}} Y_{lm}(\theta, \phi) I_{l,m} \quad (\text{outside})$$

$$\text{where } I_{l,m} = \int_0^{2\pi} \int_0^{\pi} (-\sin \phi' \hat{\mathbf{i}} + \cos \phi' \hat{\mathbf{j}}) Y_{lm}^*(\theta', \phi') \sin^2 \theta' d\theta' d\phi'$$

Expand the definition of the spherical harmonics to solve the integrals:

$$I_{l,m} = \sqrt{\frac{2l+1}{4\pi}} \sqrt{\frac{(l-m)!}{(l+m)!}} \int_0^{2\pi} (-\sin \phi' \hat{\mathbf{i}} + \cos \phi' \hat{\mathbf{j}}) e^{-im\phi'} d\phi' \int_0^{\pi} P_l^m(\cos \theta') \sin^2 \theta' d\theta'$$

Due to the orthogonality in the first integral, all terms vanish except  $m = 1$  and  $m = -1$ . The first integral can then be easily calculated to yield:

$$I_{l,\pm 1} = \sqrt{\frac{2l+1}{4\pi}} \sqrt{\frac{(l\mp 1)!}{(l\pm 1)!}} \pi (\pm i \hat{\mathbf{i}} + \hat{\mathbf{j}}) \int_0^{\pi} P_l^{\pm 1}(\cos \theta') \sin^2 \theta' d\theta'$$

Make the substitution  $x = \cos \theta$ ,  $dx = -\sin \theta d\theta$  and recognize

$$\sin \theta = \sqrt{1-x^2} = P_1^1(x) = -2P_1^{-1}(x)$$

$$I_{l,1} = \sqrt{\frac{2l+1}{4\pi}} \sqrt{\frac{(l-1)!}{(l+1)!}} \pi (i \hat{\mathbf{i}} + \hat{\mathbf{j}}) \int_{-1}^1 P_l^1(x) P_1^1(x) dx$$

$$I_{l,-1} = \sqrt{\frac{2l+1}{4\pi}} \sqrt{\frac{(l+1)!}{(l-1)!}} \pi (-i \hat{\mathbf{i}} + \hat{\mathbf{j}}) (-2) \int_{-1}^1 P_l^{-1}(x) P_1^{-1}(x) dx$$

Due to orthogonality, only the  $l = 1$  term is nonzero, leading to:

$$I_{1,\pm 1} = \sqrt{\frac{2\pi}{3}} (i \hat{\mathbf{i}} \pm \hat{\mathbf{j}})$$

The solution is now just the sum of the  $l = 1, m = -1$  term and the  $l = 1, m = 1$  term.

$$\mathbf{A} = \frac{\mu_0 \sigma \omega}{3} a r [Y_{1,1}(\theta, \phi) I_{1,1} + Y_{1,-1}(\theta, \phi) I_{1,-1}] \quad (\text{inside})$$

$$\mathbf{A} = \frac{\mu_0 \sigma \omega}{3} \frac{a^4}{r^2} [Y_{1,1}(\theta, \phi) I_{1,1} + Y_{1,-1}(\theta, \phi) I_{1,-1}] \quad (\text{outside})$$

We note that  $Y_{1,-1} = -Y_{1,1}^*$  and  $I_{1,-1} = -I_{1,1}^*$  so that:

$$\mathbf{A} = \frac{\mu_0 \sigma \omega}{3} a r [Y_{1,1}(\theta, \phi) I_{1,1} + Y_{1,1}^*(\theta, \phi) I_{1,1}^*] \quad (\text{inside})$$

$$\mathbf{A} = \frac{\mu_0 \sigma \omega}{3} \frac{a^4}{r^2} [Y_{1,1}(\theta, \phi) I_{1,1} + Y_{1,1}^*(\theta, \phi) I_{1,1}^*] \quad (\text{outside})$$

For complex numbers in general  $z + z^* = 2\Re(z)$  so that our final solutions become

$$\mathbf{A} = \frac{\mu_0 \sigma \omega}{3} a r \sin \theta \hat{\boldsymbol{\phi}} \quad (\text{inside})$$

$$\mathbf{A} = \frac{\mu_0 \sigma \omega}{3} \frac{a^4}{r^2} \sin \theta \hat{\boldsymbol{\phi}} \quad (\text{outside})$$

The magnetic flux density then becomes:

$$\mathbf{B} = \nabla \times \mathbf{A}$$

Expand in spherical coordinates, keeping only terms that are non-zero for this case:

$$\mathbf{B} = \hat{\mathbf{r}} \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_\phi) - \hat{\boldsymbol{\theta}} \frac{1}{r} \frac{\partial}{\partial r} (r A_\phi)$$

Inside:

$$\mathbf{B} = \frac{2}{3} a \mu_0 \sigma \omega [\cos \theta \hat{\mathbf{r}} - \sin \theta \hat{\boldsymbol{\theta}}]$$

$$\mathbf{B} = \frac{2}{3} a \mu_0 \sigma \omega \hat{\mathbf{z}}$$

Outside:

$$\mathbf{B} = \frac{\mu_0 \sigma \omega}{3} \frac{a^4}{r^3} [2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\boldsymbol{\theta}}]$$

$$\mathbf{B} = \frac{\mu_0 \sigma \omega}{3} \frac{a^4}{r^3} [3 \cos \theta \hat{\mathbf{r}} - \hat{\mathbf{z}}]$$